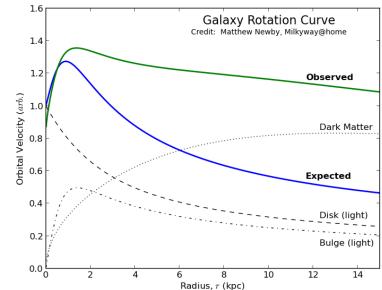
#### THE STABILITY OF CONDENSED DARK MATTER CANDIDATES

Madelyn Leembruggen University of Cincinnati Department of Physics

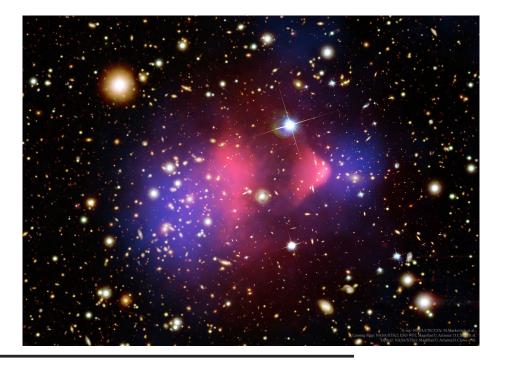


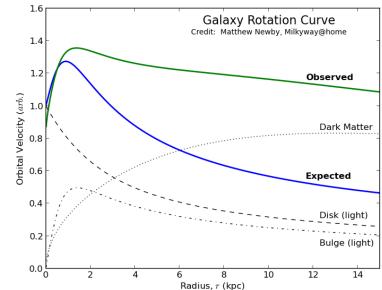
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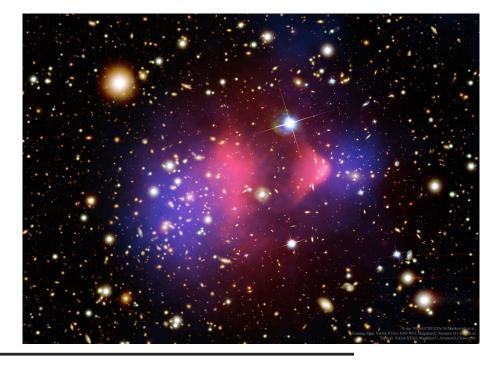


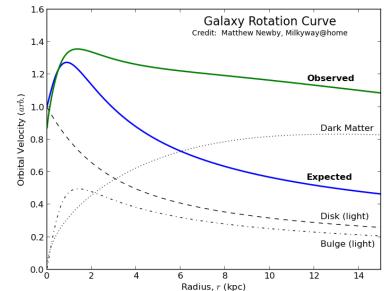
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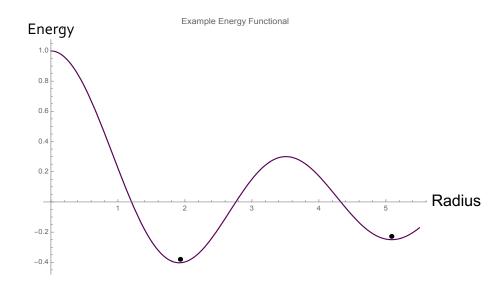
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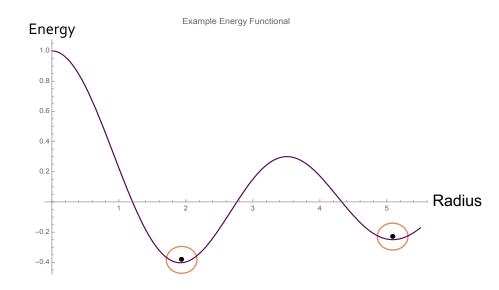


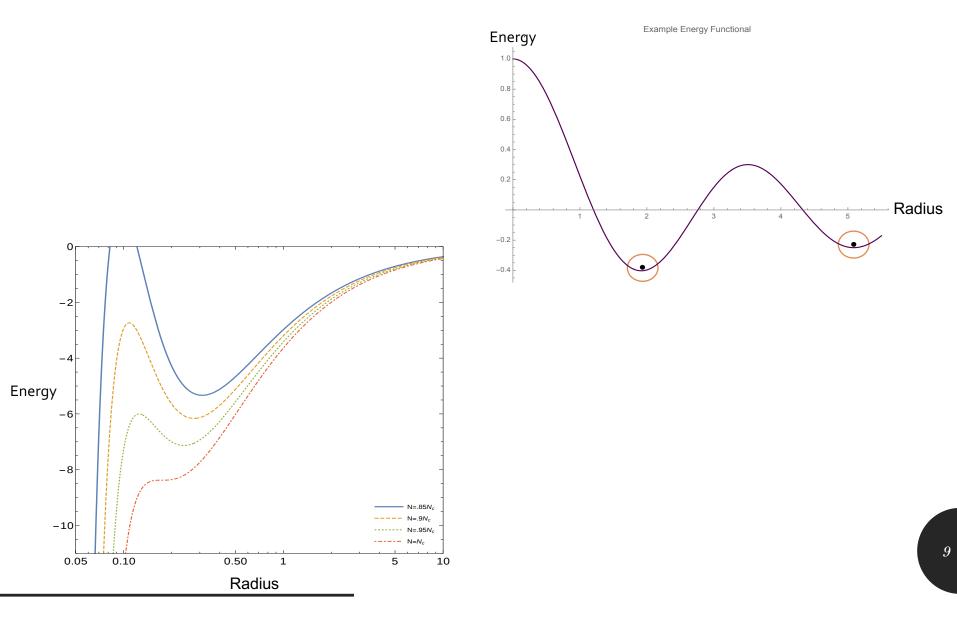




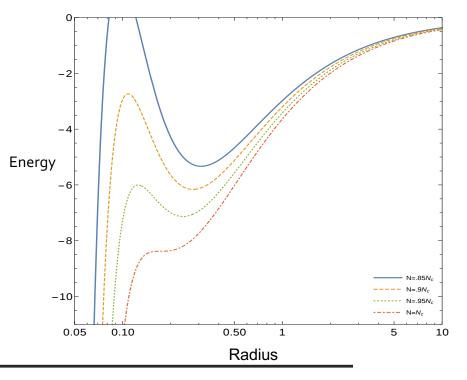
Credit: NASA/CXC/M.Weiss

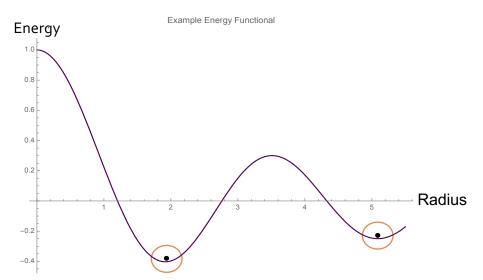




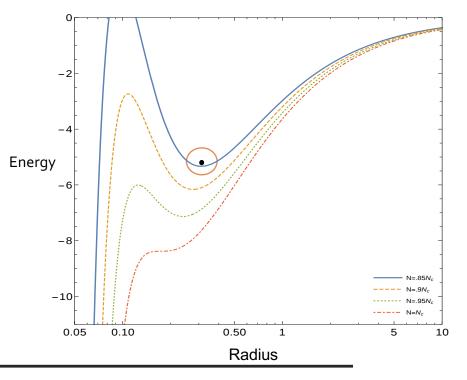


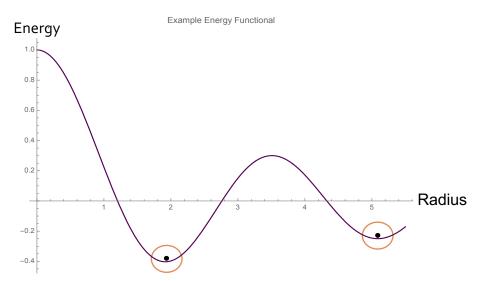
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  - Increase number of particles
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    - P. Chavanis, Phys. Rev. D. 94 (2016) 083007



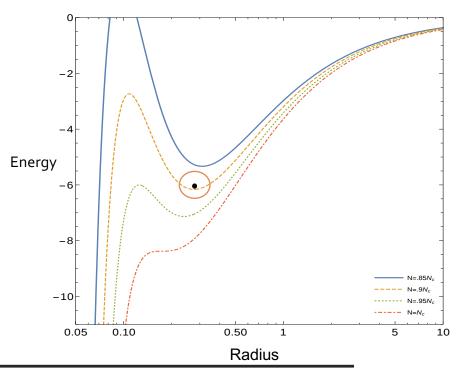


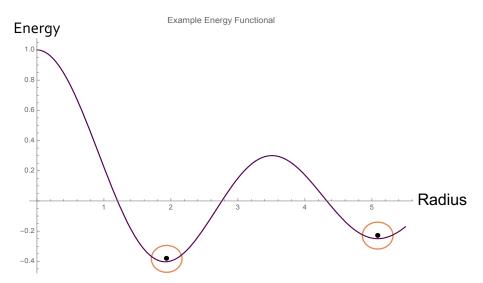
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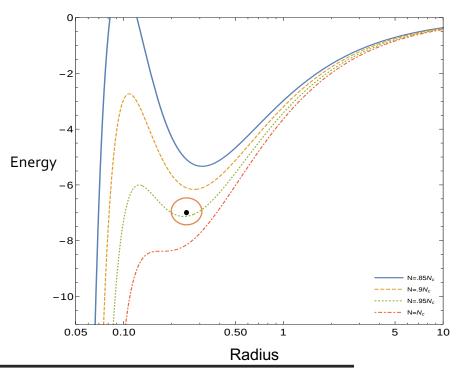


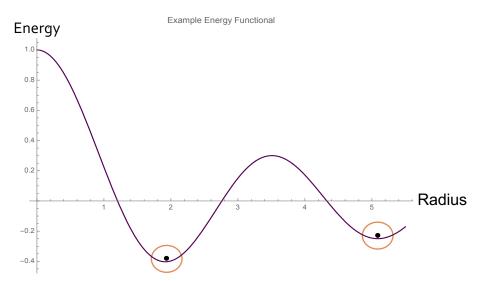
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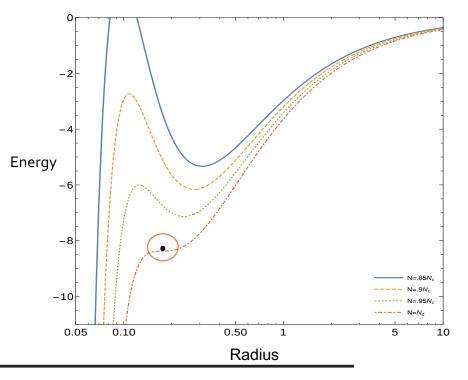


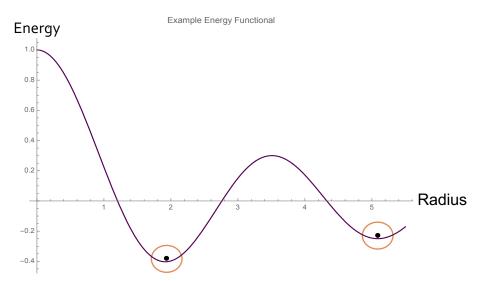
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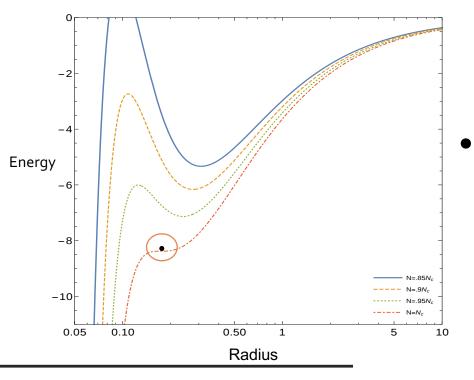


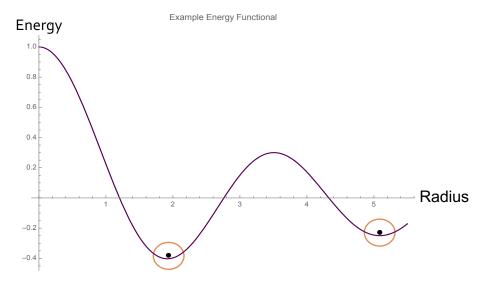
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- Stability of an axion star
  - Improve approximations
    - J. Eby, P. Suranyi, C. Vaz, L.C.R. Wijewardhana, JHEP 1503 (2015) 080 arXiv:1412.3430
  - Variational method

Axion dynamics lead us to a Hamiltonian,

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$$H = \int d^3r \left[ \frac{\left| \nabla \psi \right|^2}{2m} + W\left(\psi\right) + \frac{1}{2} V_{grav} \left| \psi \right|^2 \right]$$

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This expansion gives rise to alternating *attractive* and *repulsive* interactions!

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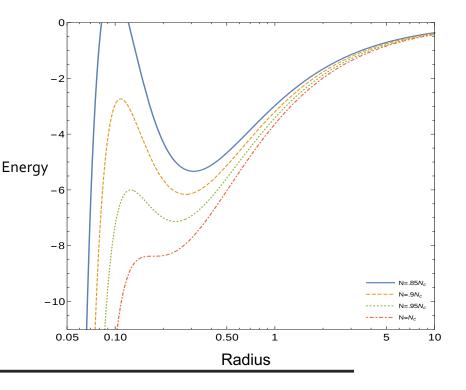
- Specify the initial conditions to match appropriate end behavior
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- So why use a variational method?
  - Numerical results exist only for stationary configurations
  - Ansätze are powerful tools for solving *dynamic* problems
  - We use numerical solutions to choose the best ansatz, but use the variational method to model collapse, collision, expansion...

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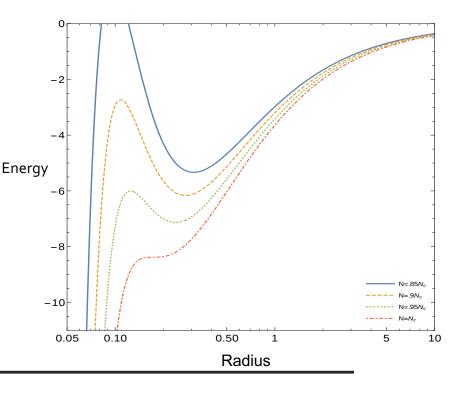
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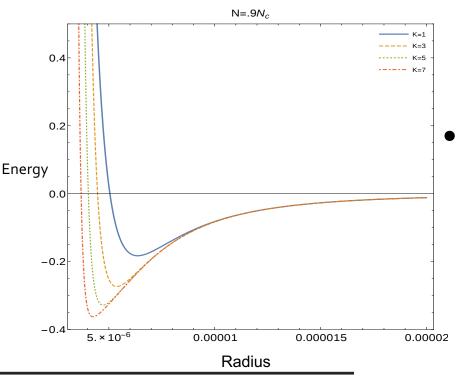
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- Approximations
  - First order approximation sufficient for large radii
  - Higher order terms necessary for small radii
  - A dense radius exists
    - Stabilized by repulsive selfinteractions
    - Corroborated by numerical solutions
    - Star will decay via a  $A_N \rightarrow A_{N-3} + a$ process

## Catalyzed Collapses

- Collisions with various astrophysical sources
  - 1. Axion star-axion star

- 2. Luminous star-axion star
- 3. Neutron star-axion star

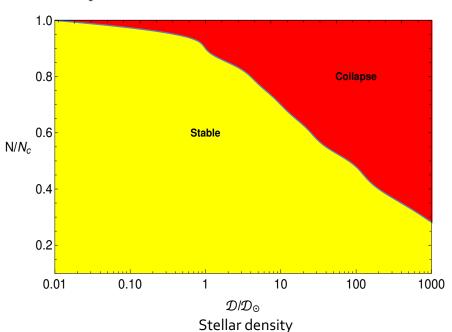
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- 10<sup>7</sup> collisions/year/galaxy, probability of collapse ~10<sup>-8</sup> per collision
- 3000 collisions/year/galaxy, collapse probability depends on stellar density
- 3. 10<sup>-3</sup> collisions/year/galaxy, axions converted to photons by B field

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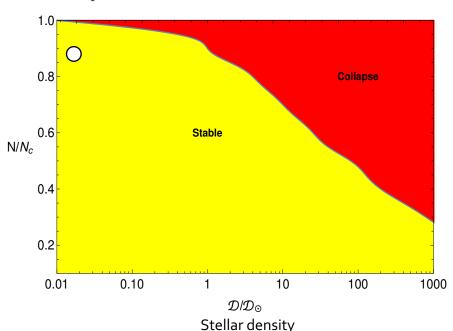


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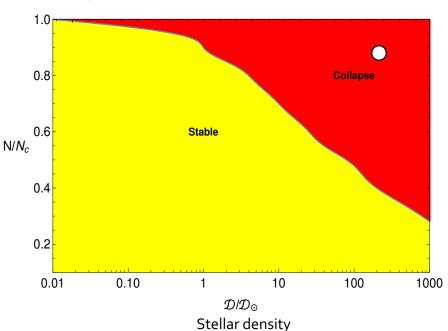


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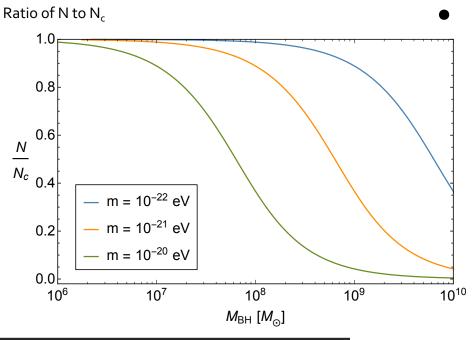
Ultra-Light Axions (ULAs)

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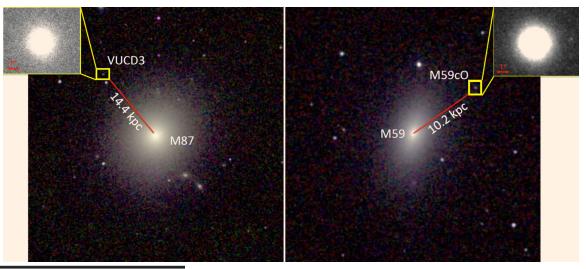
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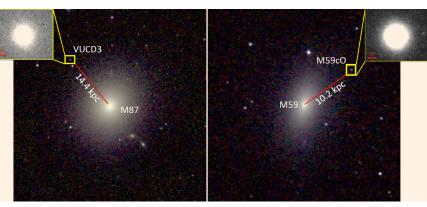
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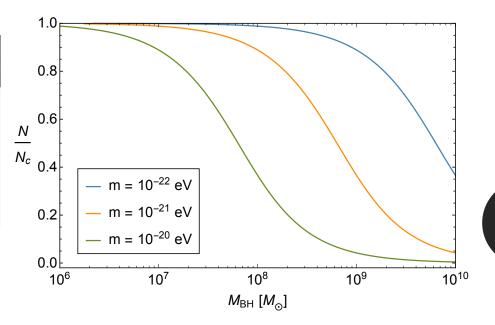


Credit: NASA/Space Telescope Science Institute

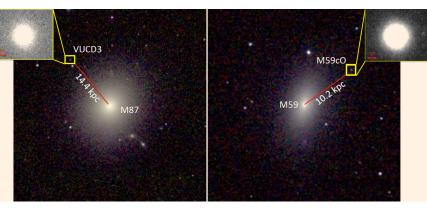
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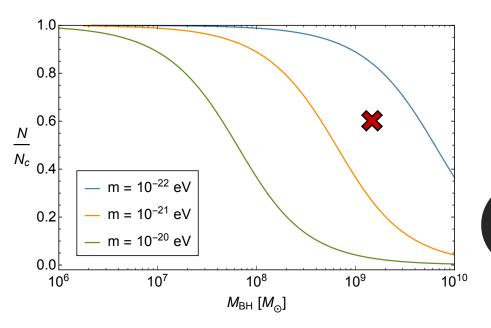
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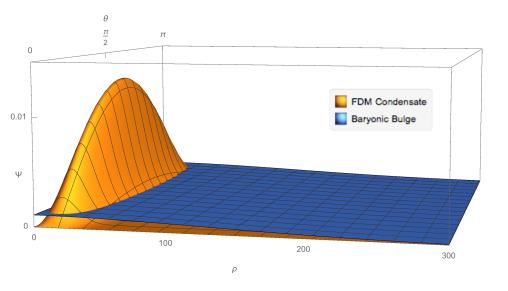
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#### <sup>10</sup> Upcoming: Rotating Dark Matter

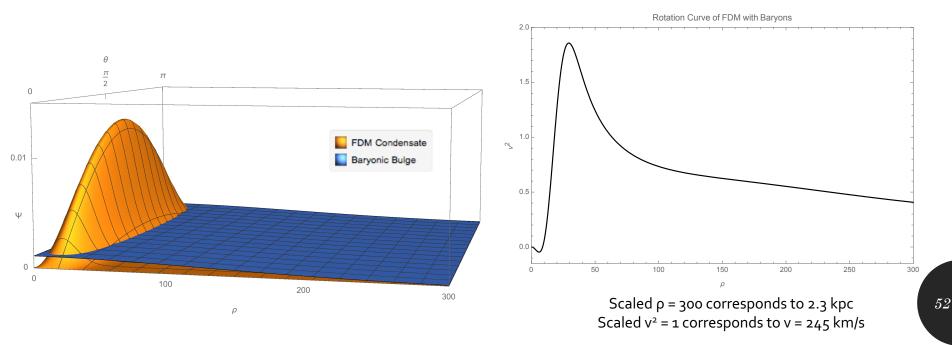
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- Galaxies rotate, shouldn't dark matter?
- Rotations affect the distribution of the dark matter BEC
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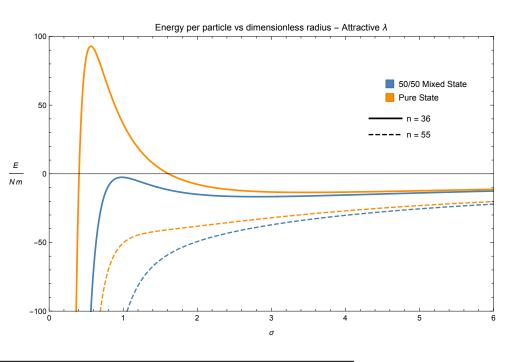


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- Is it ever energetically favorable for a self-gravitating condensate to be split between two energy levels?
  - Effective angular momentum
  - Attractive/repulsive self-interactions

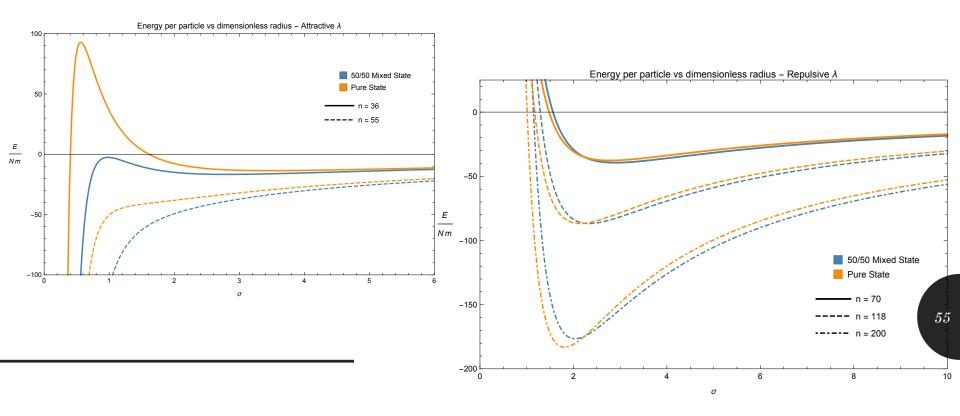
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# Concluding Thoughts

- Axions are well-motivated dark matter candidates
  - Condensates range from asteroid-size to hundreds of lightyears radii
  - Unique self-interactions
- Testable theories

- Indirect detection of QCD axion stars
- Novel use of UCD galaxies to constrain dark matter
- Formalism is broadly applicable
  - Oscillons attached to inflationary and quintessence fields
- Plenty of room to grow
  - A realm of dark photons or gauge bosons to facilitate dark matter energy transitions
  - Rotation curves to test the viability of varying bosonic candidates
  - Develop more precise methods to analyze dynamic condensates

#### Acknowledgements

- Profs. Wijewardhana, Suranyi, Ma, and Vaz
- Drs. Eby, and Gass
- L. Street and J. Leeney
- WISE Program and Prof. Ghia
- Funding from:
- The University of Cincinnati
- Department of Physics
- Barry Goldwater Scholarship Foundation

### <sup>14</sup> Papers based on this work

- Collapse of Axion Stars
  - J. Eby, M. Leembruggen, P. Suranyi, L.C.R. Wijewardhana. JHEP1612 (2016) 066. arXiv: 1608.06911
- Collisions of Dark Matter Axion Stars with Astrophysical Sources
  - J. Eby, M. Leembruggen, J. Leeney, P. Suranyi, L.C.R. Wijewardhana.
    JHEP1704 (2017) 099. arXiv: 1701.01476
- OCD Axion Star Collapse with the Chiral Potential
  - J. Eby, M. Leembruggen, P. Suranyi, L.C.R. Wijewardhana. JHEP1706 (2017) 014. arXiv: 1702.05504
- Stability of Condensed Fuzzy Dark Matter Halos
  - J. Eby, M. Leembruggen, P. Suranyi, L.C.R. Wijewardhana. Submitted to JCAP\_040P\_0718. arXiv: 1805.12147
- In preparation:
  - On Approximation Methods in the Study of Boson Stars
  - Fragmented Astrophysical Bose-Einstein Condensates

# QUESTIONS?

## <sup>16</sup> Motivation for the Axion

- The Strong CP Problem
  - QCD has very small CP violations
  - Peccei and Quinn suggest treating the CP-violation parameter as a field
    - R. D. Peccei, H. R. Quinn, Phys. Rev. Lett. 38 (1977) 1440
  - Through the "misalignment mechanism" the field produces cold axions
  - An early universe overabundance could lead to condensation
  - Spin zero bosons, electrically neutral, real scalar field
- Ultra-Light Axions
  - String theory allows for the existence of lighter classes of "axions" or axion-like particles (ALPs)
    - A. Arvanitaki et al, Phys. Rev. D 81 (2010) 123530
  - Solves cusp-core and missing satellites problems
  - Calculations can also be treated as generic bosons with various selfinteractions

#### <sup>17</sup> Searching for the Axion

• Coupled to the EM field

#### $\Delta \mathcal{L} \propto \phi \vec{E} \cdot \vec{B}$

- Added to EM Lagrangian, modifies Maxwell's equations accordingly
  - F. Wilczek, Phys. Rev. Lett. 58 (1987) 1799
- Interaction of axion field with B field produces oscillating E fields
- Frequency of E field oscillations depends on mass of axion
- Direct Detection Searches
  - ADMX, ABRACADABRA
  - CASPEr, atomic clocks
  - ALPS
  - Radio telescopes, CAST

#### Gross-Pitaevskii + Poisson from Klein-Gordon + Einstein

Klein-Gordon + Einstein equations

18

$$\Box \phi + V'(\phi) = 0 \qquad 8\pi G T_{\mu\nu} = G_{\mu\nu}$$

– Expand  $\phi$  in creation and annihilation operators:

$$\phi = ae^{-iEt}R(r) + a^{\dagger}e^{iEt}R(r) + \sum_{l\neq 0}\sum_{m} \left(a_{lm}Y_m^l R_{lm}e^{-iE_{lm}t} + h.c.\right)$$
$$|N\rangle = \frac{1}{\sqrt{N!}}a^{\dagger N}|0\rangle$$

- Assuming most of the action comes from the lowest modes, drop the higher harmonics
- Non-trivial expectation values:

$$< N |\Box \phi + V'(\phi)| N - 1 >= 0$$
  
 $< N |8\pi G T_{\mu\nu}| N >= < N |G_{\mu\nu}| N >$ 

• Write the field in terms of a wave function:

$$\phi = \frac{1}{2m} \left( \psi e^{-imt} + \psi^* e^{imt} \right)$$

#### <sup>19</sup> The Axion Potential

• Instanton Potential

$$V_I(\phi) = m^2 f^2 \left(1 - \cos\left(\frac{\phi}{f}\right)\right)$$

- $-\,$  m, f are the mass and decay parameter,  $\varphi$  is the axion field
- Chiral Potential

$$V_C(\phi) = m^2 f^2 \frac{1+z}{z} \left[ 1+z - \sqrt{1+z^2+2z\cos\left(\frac{\phi}{f}\right)} \right]$$

- $-z = m_v/m_d$ , the ratio of the up quark mass to the down quark mass
- OCD Axion Star Collapse with the Chiral Potential shows that using either potential yields the same qualitative results numerically

#### <sup>20</sup> More on Bose-Einstein Condensates

#### • BEC basics

- Fermions are subject to the Pauli Exclusion principle
- Bosons prefer the same energy state
- In low temperature/high pressure conditions, bosons condense into a macroscopic quantum system
  - Bose-Einstein Condensation in Dilute Gases by C.J. Pethick and H. Smith
- Described by a single wavefunction
- History
  - Predicted by Bose and Einstein in the 1920s
  - Created in the lab in 1995

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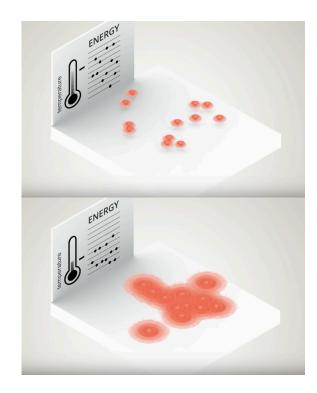
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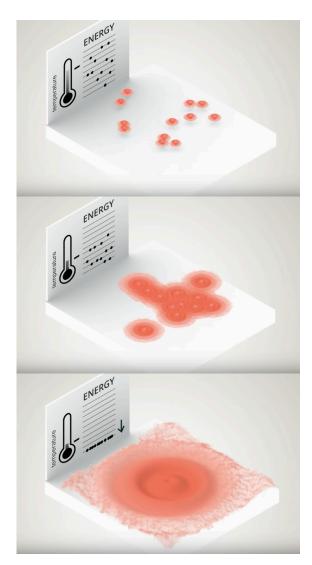
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Credit: Université Paris Sud/Labex Palm

#### More on Ultra Compact Dwarf Galaxies

- UCD formation theories
  - Small galaxies which were stripped of their stars by a larger host galaxy
  - Globular clusters from the high-mass tail of GC distribution
    - S. Mieske, M. Hilker, I. Misgeld, A&A 2012, vol. 537, A3
  - Independently formed dense galaxies
- Size and mass

- Radii in the 100 ly range
- Mass ranging up to 10<sup>10</sup> solar masses
- Often contain SMBHs
  - A.C. Seth et al, Nature 513, 398-400 (2014)

# Why doesn't the black hole eat the FDM particles?

$$\psi(r) = \sqrt{\frac{N}{\pi^{3/2}R^3}} e^{-\frac{r^2}{2R^2}} \qquad P = 4\pi \int_0^{R_{bh}} r^2 |\psi|^2 dr$$

- ULAs have deBroglie wavelength much larger than R<sub>bh</sub>
  - Somewhat analogous to optical wavelength vs lens diameter
- P(absorption by black hole) = P(particle within R<sub>bh</sub>)
- Estimate probability with integral:
  - Dilute condensate with N = 0.9  $N_c$
  - 10<sup>8</sup> solar mass black hole with corresponding Schwarzschild radius
  - $P = 5 \times 10^{-19}$