

Epsilon-Near-Zero and Plasmonic Dirac Point by using 2D materials

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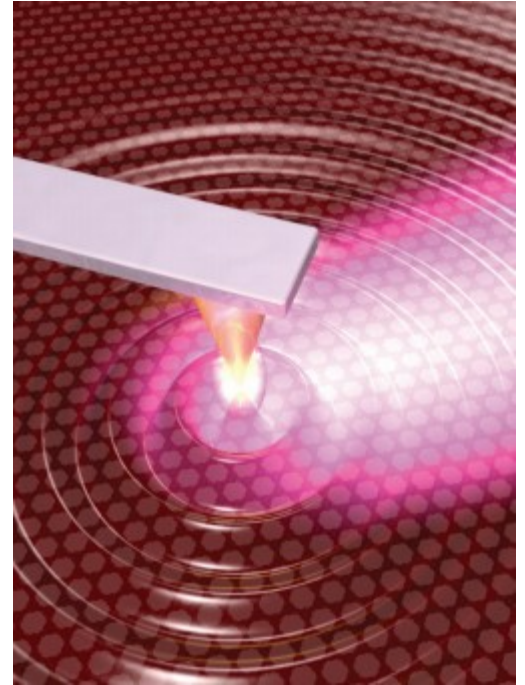
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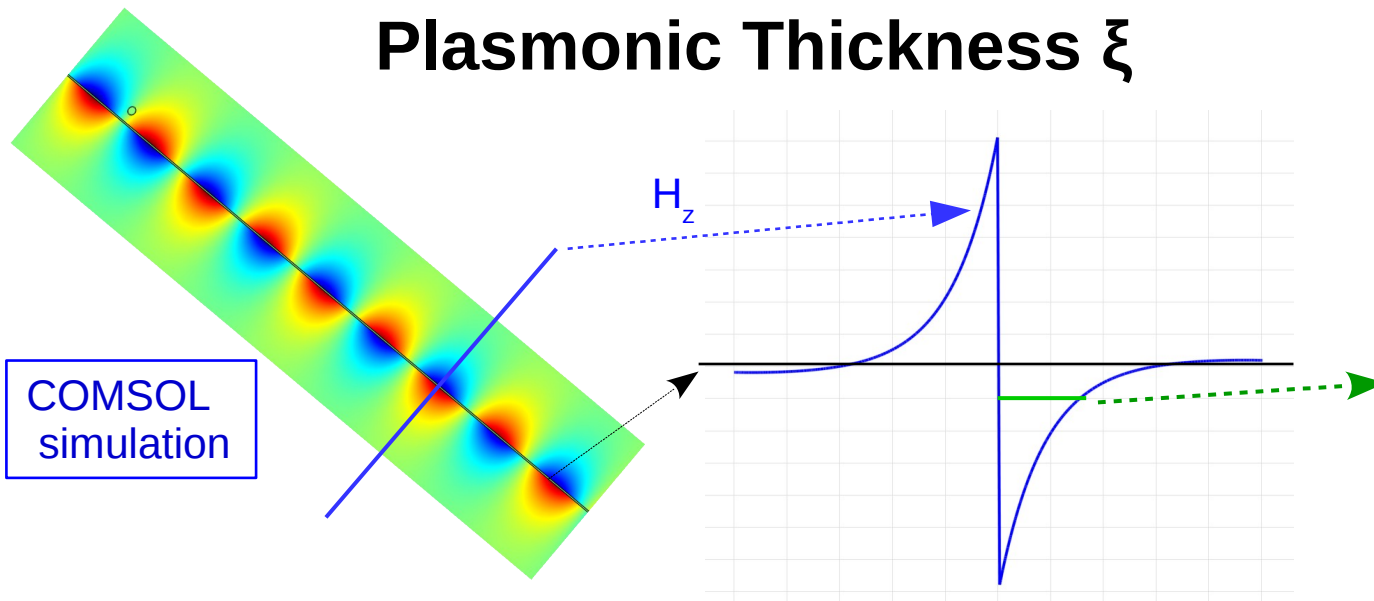
Graphene as Plasmonic Platform

Graphene is an efficient plasmonic platform:

- More confined plasmons (even more localized EM energy).
- Ultra sub-wavelength plasmons.
- Tunability of plasmon frequency via doping.
- High & low energy plasmons are supported.
- Longer propagation length.



Plasmonic Thickness ξ

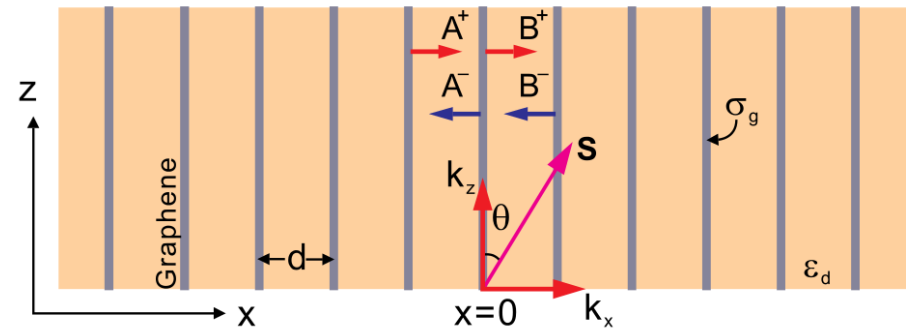
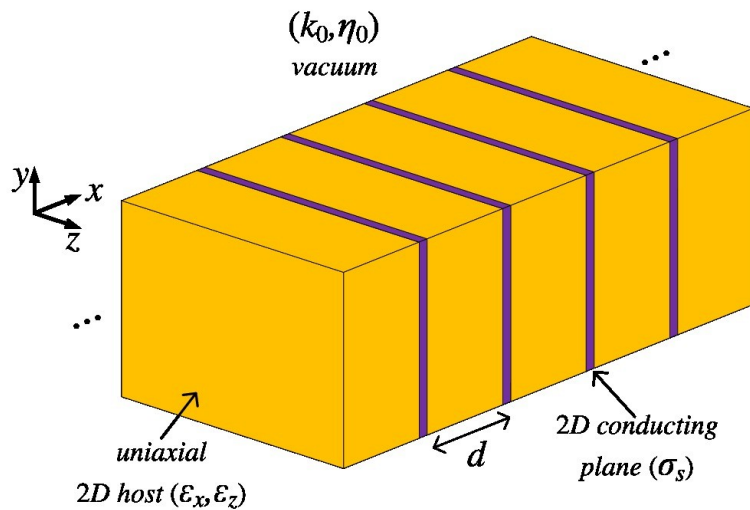


$$\delta = \xi / 2 \quad (\text{decay length})$$

$$\xi = \frac{1}{2} (k_{sp}^2 - k_0^2 \epsilon_d)^{-1/2}$$

Periodic structure of 2D metals

2D metallic layers (e.g. graphene) are extended in (y,z) plane and arranged periodically along x.



- The interlayer distance is called structural period d .
 - Anisotropic uniaxial dielectric as host media $\epsilon_z = \epsilon_y \neq \epsilon_x$.
 - The 2D layers are characterized by surface conductivity σ_s .
 - Plasmon and Bloch wavenumbers k_z and k_x .
-
- Study the normal Transverse Magnetic (TM) modes.
 - Looking for the dispersion relation: $k_z(k_x)$.
 - Due to periodicity the allowable $k_z(k_x)$ should be arranged in bands.

Maxwell Equations (MEs)

Assumptions:

- Monochromatic harmonic in time EM waves.
- Transverse Magnetic (TM) Polarization.

Maxwell Equations read:

Transversal Field

$$-i\frac{\partial}{\partial z}\begin{pmatrix} E_x \\ H_y \end{pmatrix} = k_0\eta_0 \begin{pmatrix} 0 & 1 + \frac{1}{k_0^2}\frac{\partial}{\partial x}\frac{1}{\epsilon_z}\frac{\partial}{\partial x} \\ \frac{\epsilon_x}{\eta_0} & 0 \end{pmatrix} \begin{pmatrix} E_x \\ H_y \end{pmatrix}.$$

Longitudinal Field

$$E_z = \frac{i\eta_0}{k_0\epsilon_z} \frac{\partial H_y}{\partial x}$$

$k_0 = \omega/c$ & $\eta_0 = \sqrt{\mu_0/\epsilon_0}$ are wavenumber and impedance in vacuum.

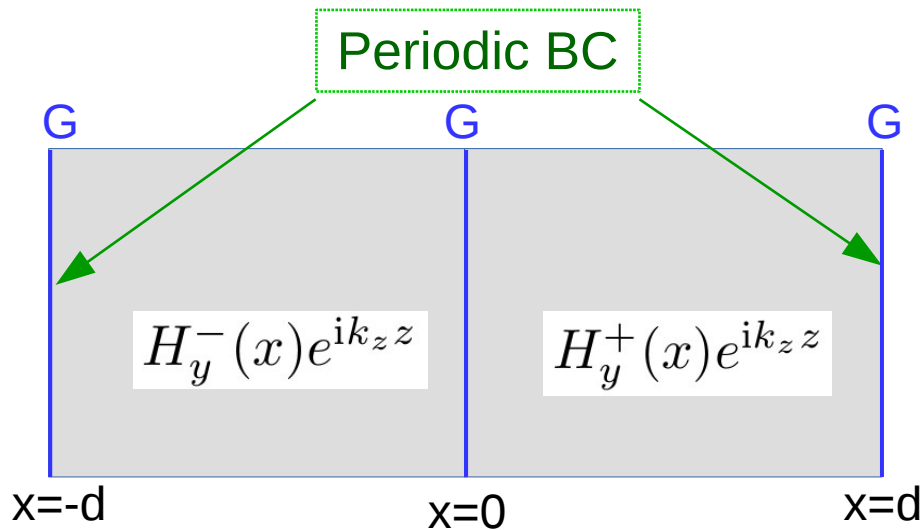
EigenValue Problem:

→ Assuming EM waves propagate along z: $\Psi(x, z) = \Psi(x)e^{ik_z z}$,

✓ Obtain an Eigenvalue problem: $k_z\Psi = \mathcal{M}\Psi$

The eigenvalue k_z is the plasmon wavenumber.

Dispersion Relation



Periodicity: The eigenmodes are Bloch waves and arranged in bands.

$$H_y^+(x) = H_y^-(x - d)e^{ik_x d}$$

$$H_y^+(0) - H_y^-(0) = \sigma_s E_z(0)$$

$$\partial_x H_y^+(0) = \partial_x H_y^-(0)$$

Graphene carries a surface current: $\mathbf{J} = \sigma \mathbf{E}_z$

$$F(k_x, k_z) = \cos(k_x d) - \left[\cosh(\kappa d) - \frac{\xi \kappa}{2} \sinh(\kappa d) \right] = 0$$

$$\kappa = \sqrt{\frac{\epsilon_z}{\epsilon_x} (k_z^2 - k_0^2 \epsilon_x)}$$

$$\xi = -\frac{i\sigma_g \eta_0}{k_0 \epsilon_z}$$

► Assume very dense grid: $\kappa d \ll 1$

► Around Brillouin center: $k_x \sim 0$.

$$\frac{k_z^2}{\epsilon_x} + \frac{d}{(d - \xi)\epsilon_z} k_x^2 = k_0^2$$

• $d > \xi$: Weak Plasmon Coupling → Elliptic Band

• $d = \xi$: Critical Plasmon Coupling → Two Linear Bands

• $d < \xi$: Strong Plasmon Coupling → Hyperbolic Band

$$k_z^2 \simeq \frac{\epsilon_x d}{\epsilon_z (\xi - d)} k_x^2$$

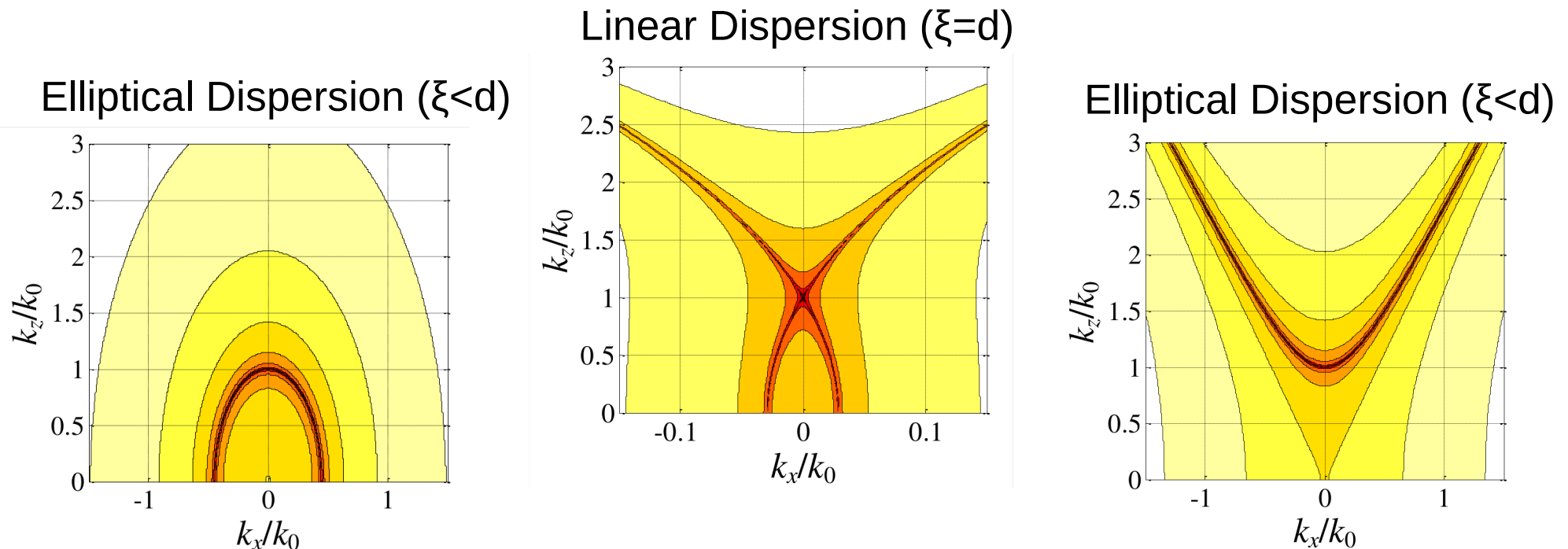
Dispersion Relation Bands

Make the choice $\xi=d$ and replace in Dispersion Relation.

- We have Saddle Point at $(k_x, k_z) = (0, k_0 \sqrt{\epsilon_x})$
- Two Bands coexist

Saddle Point + Linear Dispersion = **Plasmonic Dirac Point**

- ✓ Spatial harmonics travel with the same phase velocity.
- ✓ Non-Dispersive EM wave propagation



Effective Medium Approach

Effective Medium (Metamaterial approach)

$$\frac{k_z^2}{\varepsilon_x^{\text{eff}}} + \frac{k_x^2}{\varepsilon_z^{\text{eff}}} = k_0^2$$

Approximate Dispersion Relation

$$\frac{k_z^2}{\varepsilon_x} + \frac{d}{(d - \xi)\varepsilon_z} k_x^2 = k_0^2.$$

Effective Relative Permittivities

$$\varepsilon_x^{\text{eff}} = \varepsilon_x \quad , \quad \varepsilon_z^{\text{eff}} = \varepsilon_z + i \frac{\eta_0 \sigma_s}{k_0 d} = \varepsilon_z \frac{d - \xi}{d}.$$

Plasmonic Dirac Point → Epsilon-Near-Zero

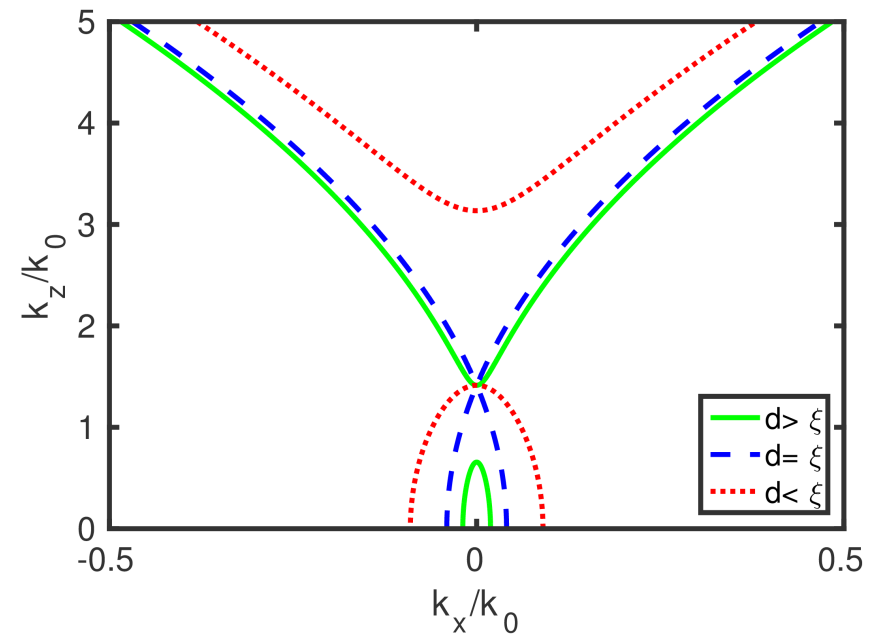
$$d = \xi \Rightarrow \varepsilon_z^{\text{eff}} = 0$$

PDP Sensitivity

PDP is very sensitive to structural defects.

$$\frac{\Delta k_z}{k_0 \sqrt{\epsilon_x}} = - \frac{6}{(k_0 d)^2 \epsilon_z} \frac{\Delta \xi}{d}.$$

$$\frac{\Delta \xi}{d} \sim \pm 0.05\%$$



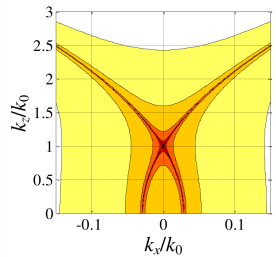
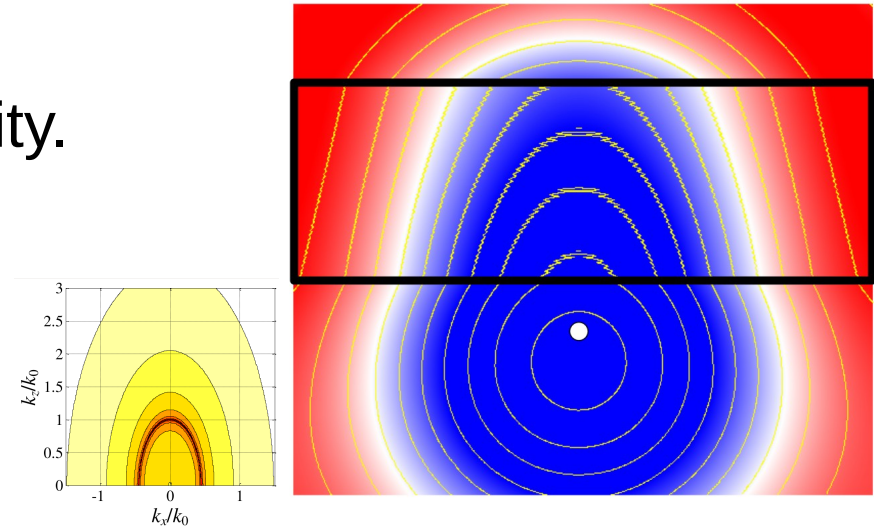
2D materials build planar bulk dielectrics.

- ✗ The extremely high sensitivity makes regular dielectrics impractical
- ✓ 2D media (e.g. MoS₂ & hBN) build bulk dielectrics with essentially perfect planarity (atomic scale controllability).

Numerical EM Wave Simulations

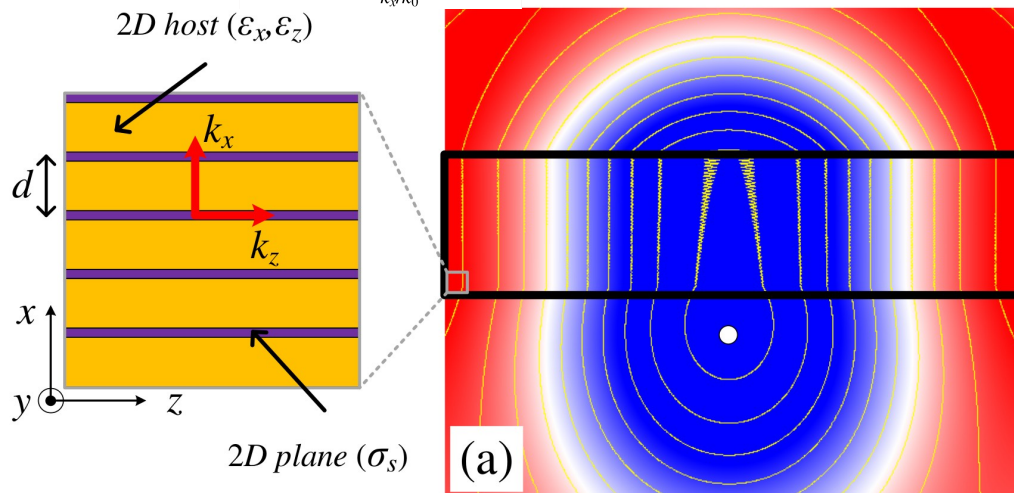
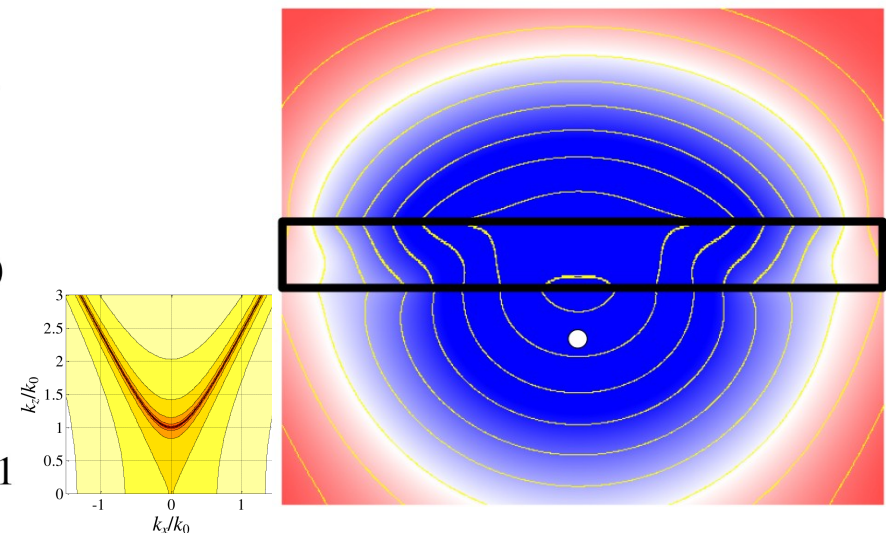
- 40 graphene monolayers embedded in MoS₂ background ($\epsilon_x=3.5$, $\epsilon_z=13$).
- Drude model for graphene conductivity.
- $\lambda_0=12 \mu\text{m}$ (THz regime).
- 2d magnetic dipole source.

Weak plasmon coupling ($\xi < d$)



PDP ($\xi=d$)

Strong plasmon coupling ($\xi > d$)



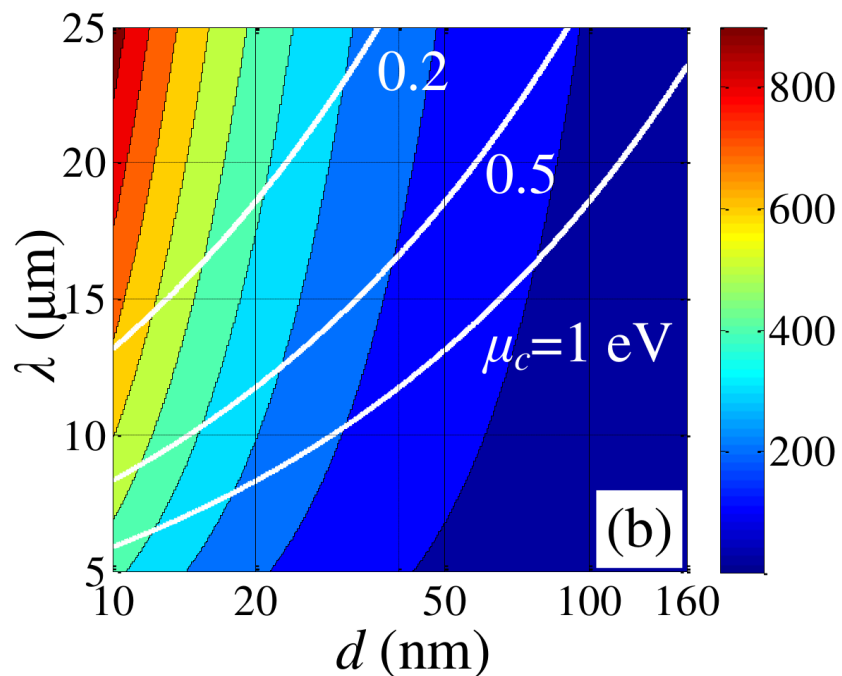
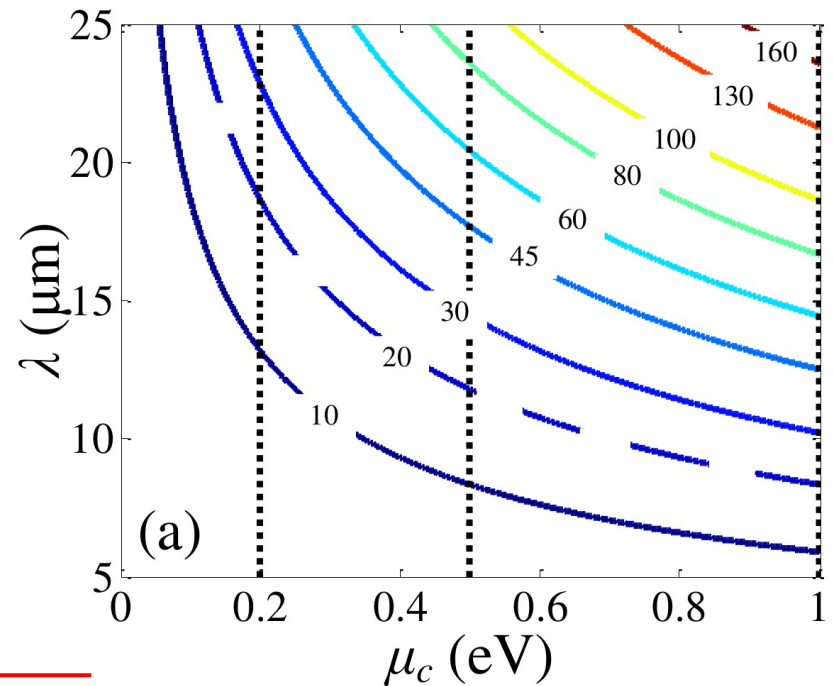
Simulations performed with COMSOL

General Investigation (maps)

Combinations of μ_c and λ leading to PDP & ENZ (ξ is plotted in nm).



A structure with arbitrary d can be fabricated and then with suitable choice of μ_c and λ we achieve ENZ.

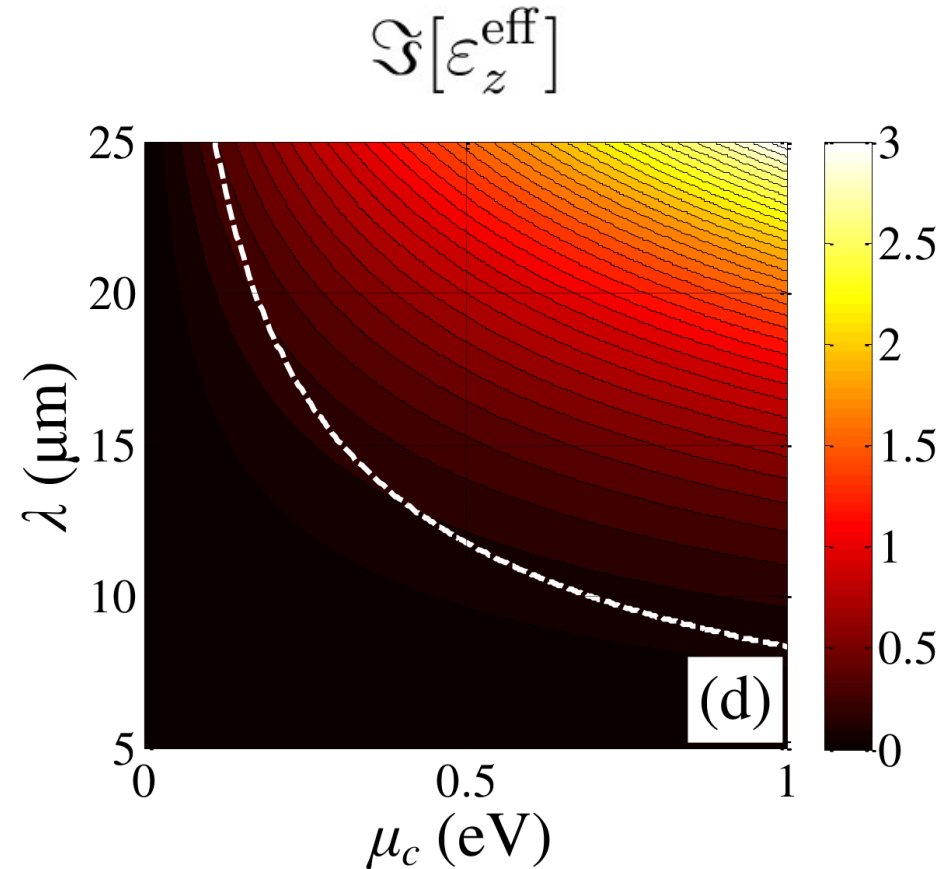
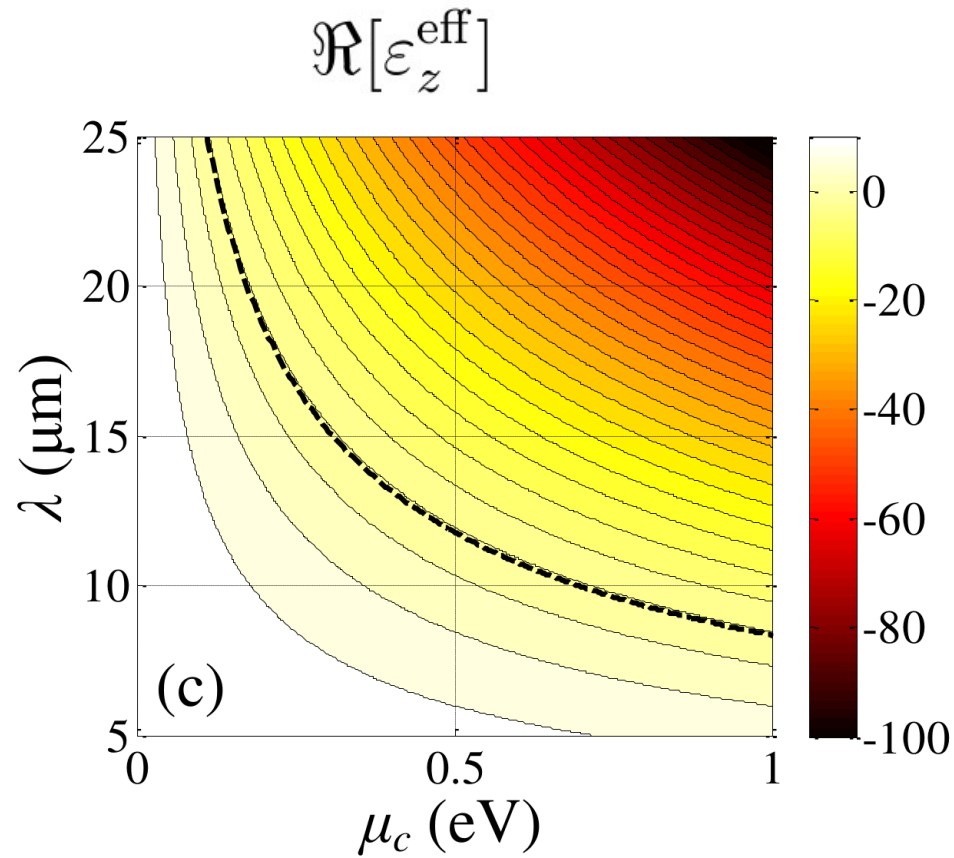


The propagation distance L/d of a plasmonic mode for all λ , d and μ_c combinations leading to ENZ.

$$\frac{L}{d} = \sqrt{\frac{2}{\epsilon_z}} \sqrt{\frac{\text{Im}[\sigma_s]}{\text{Re}[\sigma_s]}} \frac{1}{k_0 d}$$

Effective Permittivity (maps)

Effective permittivity for λ & μ_c combinations and fixed period $d=20\text{nm}$



- Dashed lines indicate the ENZ regime.
- Low losses in the ENZ region.
- Very negative ε achieved but accompanied with high losses.

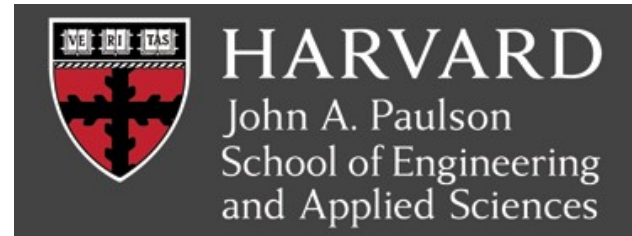
Conclusion

- Any periodic structure of 2D plasmonic materials (e.g. Graphene) exhibits Plasmonic Dirac Point in (k_x, k_z) space.
- Plasmonic Dirac Point leads to **Epsilon-Near-Zero** media.
 - Systematic method for designing ENZ meta-materials.
- Extreme sensitivity of PDP on structural imperfection.
 - Dielectrics built by 2D materials (e.g. MoS₂, hBN) have essentially perfect planarity (atomic scale control).
- Graphene multilayer in bulk MoS₂ or hBN host:
 - Dynamically tunable dispersion by doping and frequency
 - ENZ regime shows relative low losses.
 - Extremely negative (up to $\epsilon_z = -100$) relative permittivity.

Collaborators

- Prof. Efthimios Kaxiras, Harvard University.
- Prof. Giorgos Tsironis, University of Crete.
- Prof. Costas Valagiannopoulos, Nazarbayev University.
- Dr. Sharmila Shirodkar, Harvard University.

Thank you...



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